

# A Model to Predict Friction Losses of Hypoid Gears

by: H. Xu, A. Kahraman, D.R. Houser, The Ohio State University

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American Gear Manufacturers Association



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**TECHNICAL PAPER**

# **A Model to Predict Friction Losses of Hypoid Gears**

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[The statements and opinions contained herein are those of the author and should not be construed as an official action or opinion of the American Gear Manufacturers Association.]

## **Abstract**

A model to predict friction-related mechanical efficiency losses of hypoid gear pairs is proposed, which combines a commercial available finite element based gear contact analysis model and a friction coefficient model with a mechanical efficiency formulation. The contact analysis model is used to provide contact pressures and other contact parameters required by the friction coefficient model. The instantaneous friction coefficient is computed by using a validated new formula that is developed based on a thermal elastohydrodynamic lubrication (EHL) model considering non-Newtonian fluid. Computed friction coefficient distributions are then used to calculate the friction forces and the resultant instantaneous mechanical efficiency losses of the hypoid gear pair at a given mesh angle. The model is applied to study the influence of speed, load, surface roughness, and lubricant temperature as well as assembly errors on the mechanical efficiency of an example face-hobbed hypoid gear pair.

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# A Model to Predict Friction Losses of Hypoid Gears

**Hai Xu**  
Graduate Research Associate

**Ahmet Kahraman**  
Assoc. Professor

**Donald R. Houser**  
Professor Emeritus

Department of Mechanical Engineering  
The Ohio State University  
650 Ackerman Road, OH 43202

## 1 Introduction

Gear mesh friction has attracted a number of researchers for more than a century [1]. The friction between gear teeth plays an important role in defining the efficiency of the system as well as influencing scoring limits and the dynamical behavior including vibration and noise [2,3]. Both sliding and rolling actions at the gear mesh contact contribute to gear mesh friction. Sliding friction is a direct product of the relative sliding between the two contacting surfaces while rolling friction originates from the resistance to the rolling motion [4]. Coefficient of friction  $\mu$  that is used widely in the literature usually refers to the coefficient of sliding friction.

A significant number of studies have been published especially within the last forty years on friction and efficiency of gear trains as reviewed by references [5-7]. The first group of studies focused on measuring power losses of gear pair directly [8-17]. Several others measured  $\mu$  using twin-disk test machines under conditions simulating a gear pair so that this friction coefficient can be used to predict the efficiency of a gear pair [18-30, 38-39]. Some of these studies [18-25] resulted in well-known and widely used empirical formulae for  $\mu$ . These empirical formulae indicate that  $\mu$  is a function of a list of parameters such as sliding and rolling velocities, radii of curvature of the surfaces in contact, load or contact pressure, surface roughness, and the lubricant viscosity.

A group of efficiency models [31-33] investigated the efficiency of a spur gear pair by assuming a uniform  $\mu$  along the entire contact surface. A tangential friction force along the sliding direction was computed by using a given constant friction coefficient  $\mu$ , and the geometric and kinematic parameters of the spur gears. As a result, the amount of reduction of torque transmitted to the driven gear was used to calculate the mechanical efficiency  $\eta$  of the gear pair. These models were useful in bringing a qualitative understanding to the role of spur gear geometry on efficiency. They fell short in terms of the definition of  $\mu$ , as a user-defined constant  $\mu$  value must be used for every

contacting point on the tooth surface. However, the published experiments on sliding/rolling contacts indicate that many parameters might influence  $\mu$  [18-25]. In addition, these studies were limited to spur gears and many complicating effects of the tooth bending and contact deformations, tooth profile modifications and manufacturing errors were not included.

Another group of efficiency models [34-37, 40] relied on published experimental  $\mu$  formulae such as those in references [18-21]. The models in this group considered spur [35-37, 40] and helical [34] gear pairs and calculated the parameters required to define  $\mu$  according to the particular empirical formula adapted. While they are potentially more accurate than the constant  $\mu$  models, their accuracy is limited to the accuracy of the empirical  $\mu$  formula used. Each empirical  $\mu$  formula typically represents a certain type of lubricant, operating temperature, speed and load ranges, and surface roughness conditions of roller specimens that might differ from those of the gear pair that is being modeled.

The models in the last group are more advanced since they use an EHL model to predict  $\mu$  instead of relying on the user or the empirical formulae [42-54]. Among them, Dowson and Higginson [47], and Martin [48] used a smooth surface EHL model to determine the surface shear stress distribution caused by the fluid film, and hence, the instantaneous friction coefficient at the contact. Adkins and Radzimovsky [49] developed a model for lightly loaded spur gears under hydrodynamic lubrication condition and assumed that the gear tooth is rigid without deflections and local deformations. Simon [50] provided an enhancement by using point contact EHL model for heavily crowned spur gears with smooth surfaces considering the elastic displacement of the surface due to fluid pressure distributions. Larsson [51] and Wang et al [52] analyzed involute spur gear lubrication by using a transient thermal-EHL model with smooth surfaces. Wu and Cheng [53] developed a friction model based on mixed-EHL contacts and applied it to calculate the frictional power losses of

spur gears. The roughness was modeled such that all the asperities have the same radius of curvature whose heights have a Gaussian distribution. Mihalidis et al [54] included the influence of the asperity contacts as well in calculating  $\mu$  and hence efficiency. These models [47-54] were successful in eliminating to a certain extent the need for prior knowledge of  $\mu$ , at the expense of significantly more computational effort. While they were relatively enhanced in EHL aspects of the problem, the applications were limited to simple spur gears with ideal load distributions and no tooth bending deformations.

A small number of efficiency studies on helical gears were found [34,40,55-58]. Literature on hypoid gear efficiency is even sparser. Buckingham [59] proposed an approximated formula for the power loss of hypoid gears, which is the sum of the losses of a spiral bevel gear and a worm gear. Naruse et al [8,10] conducted several tests on scoring and frictional losses of hypoid gears of Klingelnberg type. Coleman [60] used a simple formula to calculate hypoid gear efficiency with a constant  $\mu$  or a  $\mu$  formula with a very limited number of parameters included [61]. Smooth-surface EHL formulations were found applied to hypoid gears by Simon [62] and Jia et al [63].

### 1.1 Objectives and Scope

Efficiency losses in a gearbox are originated from several sources including gear mesh sliding and rolling friction, windage, oil churning, and bearing friction [34]. When gears are loaded, a gear contact under load experiences combined sliding and rolling, both of which result in frictional losses. The amount of sliding frictional loss is directly related to the coefficient of friction, normal tooth load and relative sliding velocity of the surfaces while the rolling friction occurs due to the deformation of the two contacting surfaces. When the contact is lubricated, rolling frictional losses are originated from the formation of the EHL film [35]. Efficiency can be improved by reducing the coefficient of friction via precision manufacturing and smoothing the contact surfaces and enhancement of lubricant properties. Existing approaches of improving efficiency are based mostly on experimental trial-and-error type procedures focusing on such parameters, while the predictive capabilities have been limited.

The main objective of this study is to develop a mechanical efficiency model for hypoid gears. The model allows an analysis of both face-hobbed and face-milled hypoid gears. The efficiency model will allow two methods of calculating  $\mu$ , i.e. published

empirical formulae and a thermal EHL formulation. The differences amongst these approaches will be described. Parametric studies will be performed to investigate the influence of several relevant parameters such as speed, load, surface roughness, lubricant temperature as well as the assembly errors on the mechanical efficiency of hypoid gears. This study is focused primarily on the mechanical efficiency losses related to tooth friction, including sliding and rolling friction, while it relies on the published studies in terms of losses associated with windage, oil churning and bearings [34,35,64-69] when necessary.

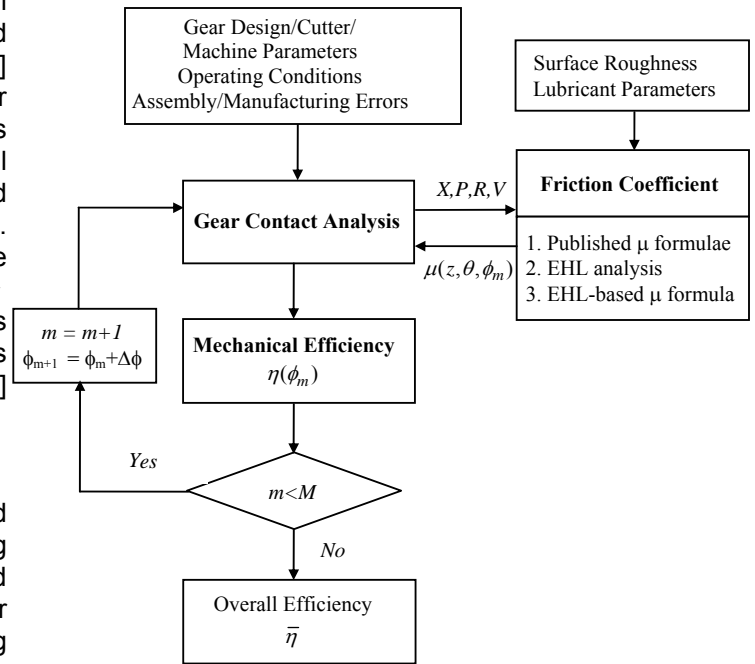


Figure 1. Flowchart for the efficiency prediction.

## 2 Efficiency Model

### 2.1 Technical Approach

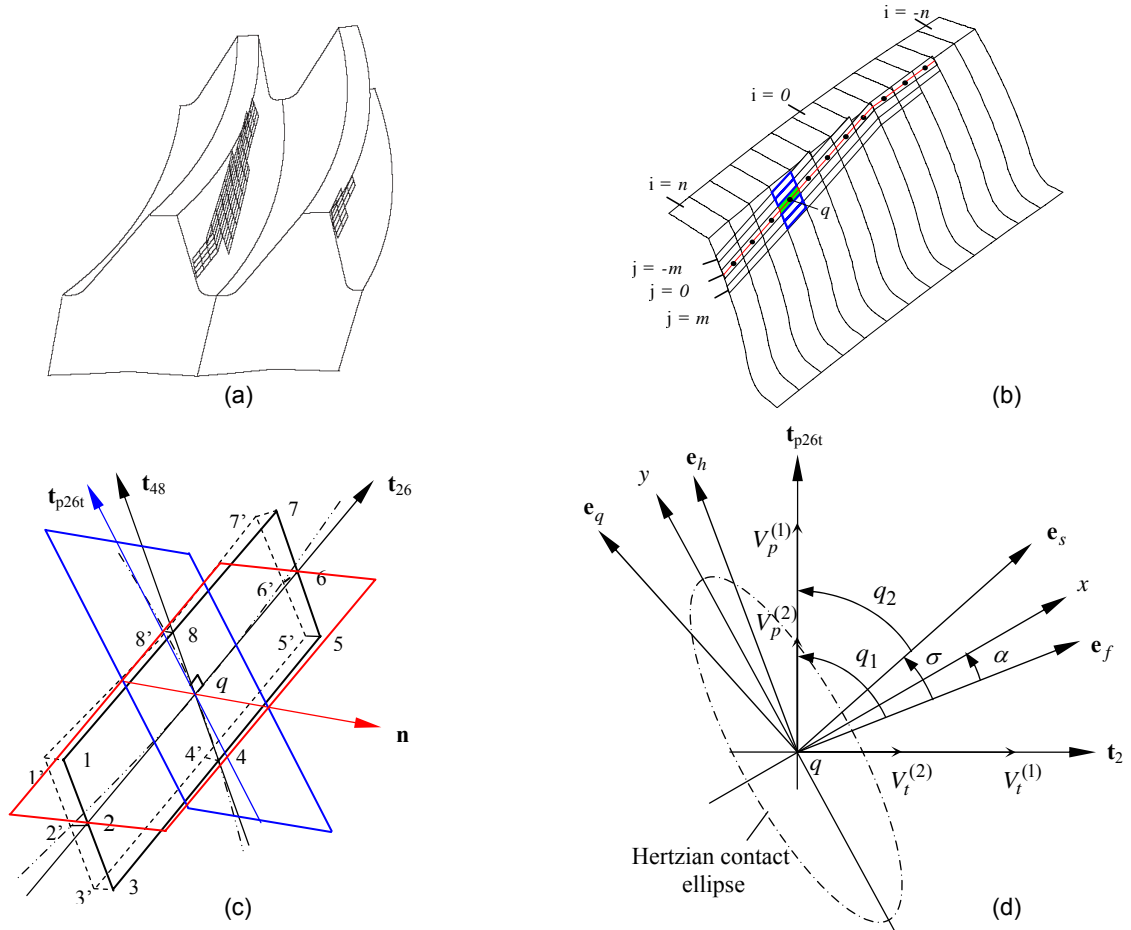
Figure 1 illustrates a flowchart of the efficiency computation methodology used in this study. Three main components are the gear contact analysis model, the friction coefficient computation model, and the gear pair mechanical efficiency computation formulation. The same methodology was applied by these authors earlier to spur and helical gears successfully [70]. It was also shown the parallel-axis efficiency model compares well with the gear pair efficiency experiments [77]. The gear contact analysis model uses the gear design parameters, operating conditions and errors associated with assembly, mounting and manufacturing of the tooth profile to predict load and contact pressure distributions at every contact point during each mesh position. Predicted load distribution or contact

pressure together with other geometric and kinematic parameters are input to the friction coefficient model to determine the instantaneous friction coefficient  $\mu(z, \theta, \phi_m)$  of every contact point  $(z, \theta)$  on the gear tooth surface.  $\mu(z, \theta, \phi_m)$  is then used by the mechanical efficiency computation module to determine the instantaneous efficiency  $\eta(\phi_m)$  of the gear pair at the  $m$ -th incremental rotational position defined by angle  $\phi_m$ . The above sequential procedure is repeated for an  $M$  number of discrete positions ( $m = 1, 2, \dots, M$ ) spaced at an increment of  $\Delta\phi$  ( $\phi_m = m \Delta\phi$ ) to cover an entire mesh cycle. These instantaneous mechanical efficiency values  $\eta(\phi_m)$  are then averaged over a complete mesh cycle to obtain the average mechanical efficiency loss of the gear pair due to tooth friction. In the following sections, main components of this methodology as shown in Fig. 1 are described in detail.

### 2.2 Contact Analysis of Hypoid Gears

A commercial available finite element (FE) based hypoid gear analysis package CALYX [71] is

used as the contact analysis tool. Both face-hobbed and face-milled versions of this model are available. The model combines FE method away from the contact zone with a surface integral formulation applied at and near the contact zone [72]. The contact analysis model used in this study has a special setup for the finite element grids inside the instantaneous contact zone. As shown in Fig. 2(a), a set of very fine contact grid is defined automatically on hypoid gear teeth to capture the entire contact zone. These grid cells are much finer than the regular size of finite element meshes elsewhere on the tooth surfaces and they are attached to the contact zones that result in more accurate contact analysis. A schematic view of these grid cells is shown in Fig. 2(b). Along the face width, there are  $2n+1$  divisions, and at each division, there is a principal contact point (shown in dot) if contact occurs. In the profile direction, there are  $2m+1$  grid cells within each division for capturing potential contact points, which would be in contact due to tooth deflections and local surface deformations.



**Figure 2.** (a) Moving grids for contact zones, (b) moving grid setup, (c) grid in tangent plane for  $\mu$  calculation, and (d) principal directions and contact ellipse.

The  $\mu$  calculation is carried out in the grid of principal contact point in the tangent plane as shown in Fig. 2(c), which is a magnified grid cell for the principal contact point  $q$ . The surface formed by dotted lines, with points 1' to 8' along the edges, is the grid on the real tooth surface and the plane formed by solid lines, with points 1 to 8 along the edges, is the grid in the tangent plane for this particular contact point  $q$ . Points 2(2'), 4(4'), 6(6') and 8(8') are the mid point at each side of the grids. Vector  $\mathbf{t}_{26}$  that connects point 2 and 6 is approximated as the instant line of contact and vector  $\mathbf{t}_{p26t}$  is normal to  $\mathbf{t}_{26}$  in the tangent plane.  $\mathbf{n}$  is the surface normal vector at the contact point  $q$ . When  $\mu$  is obtained for this principal contact point, same  $\mu$  value will be assigned to the potential contact point within the same face width division. While the load distribution and contact pressure at each grid is provided by the contact analysis model, surface velocities and curvatures are calculated in the following sections.

**Definition of Principal Contact Points.** Assume pinion surface and gear surface are defined by  $\mathbf{r}_1(s_1, t_1)$  and  $\mathbf{r}_2(s_2, t_2)$  respectively, where  $s_1, t_1$  and  $s_2, t_2$  are the surface curvilinear parameters. The principal contact point is determined and located when  $\mathbf{r}_1$  and  $\mathbf{r}_2$  become the closest to each other [72]. Surface  $\mathbf{r}_1(s_1, t_1)$  was discretized into a grid of points  $\mathbf{r}_{1ij} = \mathbf{r}_1(s_{1i}, t_{1j})$  and for each of these grid points, an effort was made to locate  $\mathbf{r}_{2ij} = \mathbf{r}_2(s_{2i}, t_{2j})$  such that  $\|\mathbf{r}_1(s_1, t_1) - \mathbf{r}_2(s_2, t_2)\|$  is minimized with respect to the variable  $s_2$  and  $t_2$ . This extremization is equivalent to solving the following system of nonlinear equations [72]

$$\begin{cases} \left[ \mathbf{r}_{1ij} - \mathbf{r}_2(s_{2i}, t_{2j}) \right] \partial \mathbf{r}_2(s_{2i}, t_{2j}) / \partial s_{2i} = 0 \\ \left[ \mathbf{r}_{1ij} - \mathbf{r}_2(s_{2i}, t_{2j}) \right] \partial \mathbf{r}_2(s_{2i}, t_{2j}) / \partial s_{2j} = 0 \end{cases} \quad (1)$$

First, solutions from Eq. (1) for each of the grid point  $\mathbf{r}_{1ij}$  were obtained by the Newton-Raphson method. Then a new grid, which is finer than the original grid, was set up around the point  $\mathbf{r}_{1ij}$  for which the separation  $\|\mathbf{r}_{1ij} - \mathbf{r}_{2ij}\|$  was the smallest. This search process is repeated several times with progressively smaller grids to locate the principal contact point [72]. The pinion and gear surfaces were divided in several parts along the face width. In each face width division, if there is contact, the principal contact is determined and then a set of grids is laid out above and below the principal contact point. These grids are added for candidate contact points that will likely be in contact due to elastic deformations of the two contacting surfaces.

**Calculation of Principal Curvatures and Principal Directions.** According to differential geometry [73], for two surfaces in contact at a common point, the principal curvatures and principal directions at the common point can be determined if the coefficients of the first and second fundamental form of the surface are given. Let  $\kappa^{(1)}$  and  $\kappa^{(2)}$  be the two eigenvalues of the eigen problem

$$\mathbf{N}\boldsymbol{\lambda} = \kappa \mathbf{M}\boldsymbol{\lambda} \quad (2)$$

where

$$\boldsymbol{\lambda} = \begin{Bmatrix} \lambda_1 \\ \lambda_2 \end{Bmatrix}, \quad \mathbf{M} = \begin{bmatrix} \frac{\partial \mathbf{r}}{\partial s} \cdot \frac{\partial \mathbf{r}}{\partial s} & \frac{\partial \mathbf{r}}{\partial s} \cdot \frac{\partial \mathbf{r}}{\partial t} \\ \frac{\partial \mathbf{r}}{\partial t} \cdot \frac{\partial \mathbf{r}}{\partial s} & \frac{\partial \mathbf{r}}{\partial t} \cdot \frac{\partial \mathbf{r}}{\partial t} \end{bmatrix},$$

$$\mathbf{N} = \begin{bmatrix} \mathbf{n} \cdot \frac{\partial^2 \mathbf{r}}{\partial s^2} & \mathbf{n} \cdot \frac{\partial^2 \mathbf{r}}{\partial s \partial t} \\ \mathbf{n} \cdot \frac{\partial^2 \mathbf{r}}{\partial t \partial s} & \mathbf{n} \cdot \frac{\partial^2 \mathbf{r}}{\partial t^2} \end{bmatrix}, \quad \mathbf{n} = \frac{\frac{\partial \mathbf{r}}{\partial s} \times \frac{\partial \mathbf{r}}{\partial t}}{\left\| \frac{\partial \mathbf{r}}{\partial s} \times \frac{\partial \mathbf{r}}{\partial t} \right\|}.$$

Here,  $\mathbf{n}$  is the unit normal vector, and  $s$  and  $t$  are the surface curvilinear parameters at the common contact point. Matrices  $\mathbf{M}$  and  $\mathbf{N}$  contain the coefficients of the first and second fundamental form of the surface, respectively. Then  $\kappa^{(1)}$  and  $\kappa^{(2)}$  are the principal normal curvatures. If the two corresponding eigenvectors  $\boldsymbol{\lambda}^{(1)}$  and  $\boldsymbol{\lambda}^{(2)}$  are normalized such that  $\{\boldsymbol{\lambda}^{(i)}\}^T \mathbf{M} \boldsymbol{\lambda}^{(i)} = 1$ ,  $i = 1, 2$ , then the two unit vectors in the principal directions corresponding to the principal curvatures are defined as [72]

$$\mathbf{e}^{(1)} = \{\boldsymbol{\lambda}^{(1)}\}^T \begin{Bmatrix} \partial \mathbf{r} / \partial s \\ \partial \mathbf{r} / \partial t \end{Bmatrix}, \quad (3a)$$

$$\mathbf{e}^{(2)} = \{\boldsymbol{\lambda}^{(2)}\}^T \begin{Bmatrix} \partial \mathbf{r} / \partial s \\ \partial \mathbf{r} / \partial t \end{Bmatrix}. \quad (3b)$$

**Determination of The Orientation of the Contact Ellipse.** Assume  $q$  is the contact point of surface  $\Sigma_1$  and surface  $\Sigma_2$ . Let  $\mathbf{e}_f$  and  $\mathbf{e}_h$  be unit vectors of principal directions, and  $\kappa_f$  and  $\kappa_h$  be principal curvatures of surface  $\Sigma_1$  at point  $q$ . Let  $\mathbf{e}_s$  and  $\mathbf{e}_q$  be unit vectors of principal directions, and  $\kappa_s$  and  $\kappa_q$  be principal curvatures of surface  $\Sigma_2$  at point  $q$ . These values of principal directions and curvatures can be solved by Eq. (2). The angle  $\sigma$ , formed by  $\mathbf{e}_f$  and  $\mathbf{e}_s$ , and the angle  $\alpha$  that determines the orientation of the coordinate axes  $x$  and  $y$  with respect to  $\mathbf{e}_f$  as shown in Fig. 2(d) can be obtained as [72]

$$\sigma = \tan^{-1} \left( \frac{\mathbf{e}_s \cdot \mathbf{e}_h}{\mathbf{e}_s \cdot \mathbf{e}_f} \right), \quad (4a)$$

$$\alpha = \frac{1}{2} \tan^{-1} \left( \frac{g_2 \sin 2\sigma}{g_1 - g_2 \cos 2\sigma} \right) \quad (4b)$$

where  $g_1 = \kappa_f - \kappa_h$ , and  $g_2 = \kappa_s - \kappa_q$ . Then the two directions  $x$  and  $y$  can be represented as

$$x = \mathbf{e}_f \cos \alpha + \mathbf{e}_h \sin \alpha, \quad (5a)$$

$$y = -\mathbf{e}_f \sin \alpha + \mathbf{e}_h \cos \alpha. \quad (5b)$$

Calculation of Surface Velocities and Radii of Curvatures in the Desired Directions. Consider that two bodies are in point contact or line contact in space and the two contacting surfaces,  $\Sigma_1$  and  $\Sigma_2$ , are in continuous tangency at the point of contact. The position vectors and unit normals of surface  $\Sigma_1$  and surface  $\Sigma_2$  must be equal. Based on this consideration, the following two relations can be obtained [74]

$$\mathbf{v}_r^{(2)} = \mathbf{v}_r^{(1)} + \mathbf{v}^{(12)} \quad (6a)$$

$$\dot{\mathbf{n}}_r^{(2)} = \dot{\mathbf{n}}_r^{(1)} + \boldsymbol{\omega}^{(12)} \times \mathbf{n} \quad (6b)$$

where  $\mathbf{v}_r^{(i)}$  is the relative velocity vector of the contact point respect to the surface as it moves over the surface  $\Sigma_i$ ,  $\dot{\mathbf{n}}_r^{(i)}$  is the velocity of the tip of the surface unit normal in its motion over the surface  $\Sigma_i$ ,  $\mathbf{n}$  is the surface unit normal vector, and  $\boldsymbol{\omega}^{(12)} = \boldsymbol{\omega}^{(1)} - \boldsymbol{\omega}^{(2)}$ . Equations (6a,b) are then used for the derivation of curvature relations of mating surfaces. As shown in Fig. 2(d), velocity vectors of point  $q$  on surface  $\Sigma_1$  can be represented in coordinate system  $S_a(\mathbf{e}_f, \mathbf{e}_h)$  as

$$\mathbf{v}_r^{(1)} = \begin{bmatrix} v_f^{(1)} & v_h^{(1)} \end{bmatrix}^T, \quad (7a)$$

$$\dot{\mathbf{n}}_r^{(1)} = \begin{bmatrix} \dot{n}_f^{(1)} & \dot{n}_h^{(1)} \end{bmatrix}^T. \quad (7b)$$

Velocity vectors of point  $q$  on surface  $\Sigma_2$  can be represented in coordinate system  $S_b(\mathbf{e}_s, \mathbf{e}_q)$  as

$$\mathbf{v}_r^{(2)} = \begin{bmatrix} v_s^{(2)} & v_q^{(2)} \end{bmatrix}^T, \quad (8a)$$

$$\dot{\mathbf{n}}_r^{(2)} = \begin{bmatrix} \dot{n}_s^{(2)} & \dot{n}_q^{(2)} \end{bmatrix}^T. \quad (8b)$$

According to Rodrigues' formula, vectors  $\mathbf{v}_r^{(i)}$  and  $\dot{\mathbf{n}}_r^{(i)}$  are collinear for the principal directions, and they are related to the principal directions as  $\dot{\mathbf{n}}_r^{(i)} = -\kappa_{I,II}^{(i)} \mathbf{v}_r^{(i)}$  [74], where  $\kappa_{I,II}^{(i)}$  are the principal

curvatures of surface  $\Sigma_i$ ,  $i=1,2$ . These principal surface are  $\kappa_f$  and  $\kappa_h$  for surface  $\Sigma_1$  and  $\kappa_s$  and  $\kappa_q$  for surface  $\Sigma_2$ . With the aid of Rodrigues' formula, Eqs. (7,8) are rearranged together with appropriate coordinate transformations to get a system of linear equations as [74]

$$\begin{bmatrix} b_{11} & b_{12} & 0 & 0 \\ b_{21} & b_{22} & 0 & 0 \\ 0 & 0 & b_{33} & b_{34} \\ 0 & 0 & b_{43} & b_{44} \end{bmatrix} \begin{Bmatrix} v_h^{(2)} \\ v_f^{(2)} \\ v_s^{(1)} \\ v_q^{(1)} \end{Bmatrix} = \begin{bmatrix} b_{15} \\ b_{25} \\ b_{35} \\ b_{45} \end{bmatrix} \quad (9)$$

where

$$\begin{aligned} b_{11} &= -\kappa_f + \kappa_s \cos^2 \sigma + \kappa_q \sin^2 \sigma, \\ b_{12} &= b_{21} = \kappa_s \sin \sigma \cos \sigma - \kappa_q \sin \sigma \cos \sigma, \\ b_{22} &= -\kappa_h + \kappa_s \sin^2 \sigma + \kappa_q \cos^2 \sigma, \\ b_{33} &= \kappa_s - \kappa_f \cos^2 \sigma - \kappa_h \sin^2 \sigma, \\ b_{34} &= b_{43} = \kappa_f \sin \sigma \cos \sigma - \kappa_h \sin \sigma \cos \sigma, \\ b_{44} &= \kappa_q - \kappa_f \sin^2 \sigma - \kappa_h \cos^2 \sigma, \\ b_{15} &= -(\boldsymbol{\omega}^{(12)} \cdot \mathbf{e}_h) - \kappa_f (\mathbf{v}^{(12)} \cdot \mathbf{e}_f), \\ b_{25} &= (\boldsymbol{\omega}^{(12)} \cdot \mathbf{e}_f) - \kappa_h (\mathbf{v}^{(12)} \cdot \mathbf{e}_h), \\ b_{35} &= -(\boldsymbol{\omega}^{(12)} \cdot \mathbf{e}_q) - \kappa_s (\mathbf{v}^{(12)} \cdot \mathbf{e}_s), \\ b_{45} &= (\boldsymbol{\omega}^{(12)} \cdot \mathbf{e}_s) - \kappa_q (\mathbf{v}^{(12)} \cdot \mathbf{e}_q). \end{aligned}$$

Equation (9) can then be solved to obtain the surface velocities in the principal directions. The surface velocities in the direction of contact line and in the direction that is normal to the contact line for surface  $\Sigma_1$  and surface  $\Sigma_2$ , as shown in Fig. 2(d), can then be obtained by coordinate transformation as follows

$$\begin{bmatrix} V_p^{(1)} \\ V_t^{(1)} \end{bmatrix} = \begin{bmatrix} \cos q_2 & \sin q_2 \\ -\sin q_2 & \cos q_2 \end{bmatrix} \begin{bmatrix} v_s^{(1)} \\ v_q^{(1)} \end{bmatrix} \quad (10a)$$

$$\begin{bmatrix} V_p^{(2)} \\ V_t^{(2)} \end{bmatrix} = \begin{bmatrix} \cos q_1 & \sin q_1 \\ -\sin q_1 & \cos q_1 \end{bmatrix} \begin{bmatrix} v_f^{(2)} \\ v_h^{(2)} \end{bmatrix} \quad (10b)$$

Then, sliding and rolling velocities in the direction of the contact line can be obtained as  $V_{st} = V_t^{(1)} - V_t^{(2)}$  and  $V_{rt} = V_t^{(1)} + V_t^{(2)}$ , respectively. Sliding and rolling velocities in the direction that is normal to the contact line are  $V_{sp} = V_p^{(1)} - V_p^{(2)}$  and  $V_{rp} = V_p^{(1)} + V_p^{(2)}$ , respectively. The resultant sliding and rolling velocities are given as  $V_{stotal} = (V_{st}^2 + V_{sp}^2)^{0.5}$  and  $V_{rtotal} = (V_{rt}^2 + V_{rp}^2)^{0.5}$ , respectively. Normal curvatures in the direction of the contact line,  $\kappa_{nc}^{(i)}$ , and in the direction that is normal to the contact line,  $\kappa_{nm}^{(i)}$ , can be obtained as follows

$$\kappa_{nc}^{(i)} = \kappa_I^{(i)} \cos^2(\pi/2 + q_i) + \kappa_{II}^{(i)} \sin^2(\pi/2 + q_i) \quad (11a)$$

$$\kappa_{nm}^{(i)} = \kappa_I^{(i)} \cos^2 q_i + \kappa_{II}^{(i)} \sin^2 q_i \quad (11b)$$

**Table 1** Empirical friction coefficient formulae considered in this study ( $\nu_k$  and  $\nu$  are given in centistokes and centipoises, respectively).

Formula and author	Applicable parameter ranges	Specific units
Drozdov and Gavrikov [21]: $\mu = [0.8\sqrt{\nu_k V_s + V_r \phi} + 13.4]^{-1}$ $\phi = 0.47 - 0.13(10)^{-4} P_{\max} - 0.4(10)^{-3} \nu_k$	$\nu_k \in [4, 500]$ $V_s \leq 15$ , $V_r \in [3, 20]$ $P_{\max} \in [4000, 20000]$	$V_s, V_r = m/s$ $P_{\max} = kg/cm^2$
O'donoghue and Cameron [20] $\mu = 0.6[(S + 22)/35] \left[ \nu^{1/8} V_s^{1/3} V_r^{1/6} R^{1/2} \right]^{-1}$		$S = \mu in$ , CLA $V_s, V_r = in/s$ , $R = in$
Misharin [18] $\mu = 0.325[V_s V_r \nu_k]^{-0.25}$	$V_s/V_r \in [0.4, 1.3]$ $P \geq 2,500$ , $\mu \in [0.02, 0.08]$	$V_s, V_r = m/s$ $P = kg/cm^2$
ISO TC60 [41] $\mu = 0.12[W' S / (R V_r \nu)]^{0.25}$		$V_r = m/s$ , $R = mm$ $S = \mu m$ , RMS, $W' = N/mm$
Benedict and Kelley [19] $\mu = 0.0127 \left[ \frac{50}{50 - S} \right] \text{Log}_{10} \left[ \frac{3.17(10)^8 W'}{\nu V_s V_r^2} \right]$	$\frac{50}{50 - S} \leq 3$	$S = \mu in$ , RMS $W' = lbf/in$ $V_s, V_r = in/s$

where,  $q_i$  is the angle that is formed by the direction of interest and  $e_f$  and  $e_s$  for  $i=1$  and  $i=2$ , respectively. Then the radius of curvature is simply the inverse of the corresponding normal curvature value in Eq. (11).

### 2.3 Friction Coefficient Models

**Empirical Friction Coefficient Formulae.** A large number of empirical formulae for coefficient of friction can be found in the literature. Most of these formulae obtained from measured data of twin-disk type tests were in the following general form

$$\mu = f(\nu_k, \nu, V_s, V_r, R, W', P_{\max}, S, \dots). \quad (12)$$

Here  $\nu_k$  and  $\nu$  are kinematic and dynamic viscosity of the lubricant, both of which are measured at the oil inlet, and are function of inlet oil temperature at ambient pressure. Parameters  $V_s$ ,  $V_r$ , and  $R$  denote relative surface sliding velocity, sum of the rolling velocities and the combined radius of curvature, respectively. The load parameters are the unit normal load  $W'$  or the contact pressure  $P_{\max}$ .  $S$  is a surface finish parameter, representing the initial composite surface roughness of the two contacting surfaces.

A representative set of commonly cited  $\mu$  formulae given in Table 1 are considered in this study. These formulae are quite different from each other as they consider different parameters and parameter ranges. Formulae by Drozdov and Gavrikov [21] and Misharin [18] include neither any surface roughness parameter nor a radius of

curvature  $R$ , while others consider a surface roughness parameter  $S$  (root-mean-square (RMS) or centerline average (CLA)). ISO TC60 [41] formula does not include  $V_s$  while formulae of O'donoghue and Cameron [20] and Misharin [18] do not consider the load as a parameter and hence they should not be able to account for any load effects on  $\mu$ . In addition, each formula listed in Table 1 is valid for a range of key system parameters. For instance, the formula proposed by Drozdov and Gavrikov is valid for  $\nu_k \in [4, 500]$  centistokes,  $V_s \leq 15 m/s$ ,  $V_r \in [3, 20] m/s$  and  $P_{\max} \in [4000, 20000] kg/cm^2$ , and would be suitable for a gear application only if these gear contact parameters stay within these ranges.

**Calculation of Friction Coefficient Using an EHL Model.** According to EHL theory, distribution of the film pressure and thickness across an elastic lubricated contact can be obtained by solving the transient Reynolds equation simultaneously with a number of other equations including the film thickness equation, the energy balance equation, viscosity-pressure-temperature relationship, density-pressure-temperature relationship, and the load equation [75]. In this study, a deterministic EHL model based on the model proposed by Cioc et al [75] will be used. Assuming that a hypoid gear pair exhibits line contact conditions, a non-dimensional transient Reynolds equation is written as

$$\frac{\partial}{\partial X} \left[ s \left( \frac{P_h b}{12 \nu_0 \mu} \right) \left( \frac{\bar{\rho} H^3}{\bar{\nu}} \right) \frac{\partial P}{\partial X} \right] = \frac{\partial}{\partial X} (\bar{\rho} H) + \frac{\partial}{\partial t} (\bar{\rho} H) \quad (13)$$



with the boundary conditions of  $P=0$  at inlet, and  $P=0$  and  $\partial P/\partial X=0$  at outlet. In Eq. (13),  $b$  is the half width of Hertzian contact zone, and  $X$  and  $H$  are the dimensionless coordinate and the dimensionless film thickness, respectively, both normalized by  $b$ ,  $P_h$  is the maximum Hertzian pressure,  $\bar{\nu}$  is the dimensionless dynamic viscosity normalized by  $\nu_0$ , which is the dynamic viscosity at oil inlet, and  $\bar{\rho}$  is the dimensionless lubricant density normalized by the lubricant density at inlet.  $u$  is the rolling velocity defined as  $u=(u_1+u_2)/2$  and  $\bar{t}$  is the dimensionless time defined as  $\bar{t}=ut/b$ . For a Newtonian fluid  $s=1$  while  $s$  is defined by a nonlinear function for a non-Newtonian fluid [75]. The film thickness equation has the form

$$H = \frac{h_0}{b} + \left(\frac{b}{2R}\right)X^2 + \frac{d(X,\bar{t})}{b} + S(X,\bar{t}) \quad (14)$$

where,  $h_0$  is the reference film thickness,  $d(X,\bar{t})$  is the total elastic deformation,  $S(X,\bar{t})$  is the composite surface profile irregularities, and  $R$  is the combined radius of curvature of the contact. A two-slope viscosity-pressure-temperature model [75] is used here as

$$\nu = \begin{cases} e^{[G_1 P - \bar{\gamma}(T-1)]}, & P < P_a, \\ e^{[c_0 + c_1 P + c_2 P^2 + c_3 P^3 - \bar{\gamma}(T-1)]}, & P_a \leq P \leq P_b, \\ e^{[G_1 P + G_2(P-P_b) - \bar{\gamma}(T-1)]}, & P_b \leq P. \end{cases} \quad (15)$$

Here,  $G_1 = \alpha_1 P_h$ ,  $G_2 = \alpha_2 P_h$ ,  $\bar{\gamma}$  is the dimensionless coefficient of thermal expansion, and  $T$  is the dimensionless film temperature normalized by the oil temperature at inlet.  $P_t$  is the transitional pressure, a threshold value beyond which an increase in viscosity changes its slope from  $\alpha_1$  to  $\alpha_2$ . Transition pressure values of  $P_a = 0.7P_t$  and  $P_b = 1.4P_t$  are used with  $P_t = 380$  MPa as suggested by Allen for mineral oils [75]. The coefficients  $c_0$  to  $c_3$  are determined such that that transition between the two slopes is smooth. In addition, the following density-pressure-temperature relationship is considered

$$\bar{\rho} = \left(1 + \frac{C_a P}{1 + C_b P}\right) [1 - \bar{\beta}(T-1)] \quad (16)$$

where  $C_a = 0.6 \times 10^{-9} P_h$ ,  $C_b = 1.7 \times 10^{-9} P_h$ , and  $\bar{\beta}$  is the dimensionless coefficient of thermal expansion [75].

Since a thermal EHL model is sought, the energy equation for non-Newtonian lubricants is given as

$$\frac{\partial^2 T}{\partial Z^2} = C_1 U \bar{\rho} \frac{\partial T}{\partial X} - C_2 T U \frac{\partial P}{\partial X} - C_3 \frac{\bar{\tau}_1 \bar{\tau}_L}{\bar{\nu} \sqrt{(\bar{\tau}_L)^2 - (\bar{\tau}_1)^2}} + C_4 \bar{\rho} \frac{\partial T}{\partial \bar{t}} \quad (17)$$

In Eq. (17), coefficients  $C_i$  ( $i=1$  to  $4$ ) are functions of thermal parameters,  $Z$  is the dimensionless coordinate normalized by  $b$ ,  $U$  is the dimensionless fluid velocity,  $\bar{\tau}_1$  is the dimensionless shear stress of surface 1, and  $\bar{\tau}_L$  is the dimensionless limiting shear stress. Finally, the load distribution predicted should be such that the normal load applied at the contact is balanced, i.e.

$$\bar{W} = \int_{X_{in}}^{X_{out}} [P(X,\bar{t}) + P_c(X,\bar{t})] dX \quad (18)$$

where  $\bar{W}$  is the dimensionless unit load applied,  $X_{in}$  and  $X_{out}$  are the inlet and outlet of computational domain for the contact zone,  $P_c$  is the asperity contact pressure for where asperity contact occurs. This load balance equation is checked periodically and the reference film thickness  $h_0$  in Eq. (14) is adjusted, if necessary.

In reference [75], Eqs. (13-18) were solved numerically by using the finite difference method. Using the same methodology, distributions of the pressure  $p(x)$ , viscosity  $\nu(x)$ , and fluid film thickness  $h(x)$  across the lubricated contact zone are predicted. In the case of no asperity contacts, the surface shear stress distribution across the Hertzian contact width  $-b \leq x \leq b$  is given by

$$\tau(x) = -\frac{h(x)}{2} \frac{\partial p}{\partial x} + \nu(x) \frac{u_1 - u_2}{h(x)}. \quad (19)$$

Typical ground or shaved gear tooth surfaces cannot be considered smooth unless they are polished chemically or mechanically. The actual surface roughness amplitudes are typically comparable to the fluid film thickness. As a result, the portion of the normal load might be carried by actual asperity contacts resulting in mixed- or partial-EHL conditions. In this case, Eq. (19) is still valid at the  $x$  locations where a fluid film is maintained between the two surfaces. At the  $x$  locations where actual asperity contacts take place

$$\tau(x) = \mu_s p(x) \quad (20)$$

where  $\mu_s$  is the friction coefficient between the contacting asperities. Consequently, the friction traction force per unit width of the contact can be written as

$$F_t = \int \tau(x) dx = F_r' + F_s' \quad (21)$$

where  $F_r'$  is the rolling friction per unit width of the contact that comes from the viscous flow in the formation of the film thickness and squeeze motion during pure rolling,

$$F_r' = -\int \frac{h(x)}{2} \frac{\partial p}{\partial x} dx \quad (22)$$

and  $F_s'$  is the sliding friction per unit width of the contact that is caused by fluid shear and the asperity contact, respectively [76]

$$F_s' = \int v(x) \frac{u_1 - u_2}{h(x)} dx + \int \mu_s p(x) dx . \quad (23)$$

Finally, the coefficient of sliding friction of the entire section of contact is given by  $\mu = F_s'/W'$  where  $W'$  is the applied normal load per unit width of contact.

With the exception of  $\mu_s$  in Eq. (23) that must be defined in cases when asperity contact occurs, the approach outlined above to compute  $\mu$  is entirely physics-based. One disadvantage of this approach is that it requires significant computational effort since a large number of individual EHL analysis must be carried out for each discretized segment along the lines of contact at each rotational gear position considered.

## 2.4 Gear Efficiency Computation

Once the coefficient of friction  $\mu(z, \theta, \phi_m)$  at each contact point  $(z, \theta)$  at each rotational angle  $\phi_m$  ( $m=1, 2, \dots, M$ ) are known, the sliding friction force at each contact position can be calculated by

$$F_s(z, \theta, \phi_m) = \mu(z, \theta, \phi_m) W(z, \theta, \phi_m) \quad (24)$$

where  $W(z, \theta, \phi_m)$  is the normal load at the contact point of interest. The rolling friction force  $F_r$  can be obtained by multiplying  $F_r'$  in Eq. (22) with the width of the contact. When the empirical formulae like the ones listed in Table 1 are used,  $F_r$  can be estimated from an isothermal EHL contact [53] as

$$F_{ro} = 4.318 (\tilde{G}\tilde{U})^{0.658} \tilde{Q}^{0.0126} R / \alpha \quad (25)$$

where  $\tilde{G} = \alpha E'$  is the dimensionless material parameter,  $\tilde{U} = v_0(u_1 + u_2)/(E'R)$  is the dimensionless speed parameter,  $\tilde{Q} = W'/(E'R)$  is the dimensionless load parameter,  $R$  is the effective radius of curvature,  $\alpha$  is the pressure viscosity coefficient. In order to account for the effect of temperature rise at high speed conditions, a thermal reduction factor  $\varphi_T$

can be used to modify this isothermal formula such that

$$F_r = \varphi_T F_{ro} \quad (26)$$

where  $\varphi_T$  is defined in reference [53] as

$$\varphi_T = \frac{1 - 13.2(P_h/E')(L^*)^{0.42}}{1 + 0.213(1 + 2.23SR^{0.83})(L^*)^{0.64}} , \quad (27)$$

$$L^* = -\frac{\partial v}{\partial t_0} \frac{(V_e)^2}{K_f} . \quad (28)$$

Here,  $v$  is the absolute viscosity in cPs,  $t_0$  is the lubricant temperature at inlet in degrees C,  $K_f$  is the lubricant thermal conductivity in  $W/m^\circ C$ ,  $SR$  is the slide-to-roll ratio that is defined as  $SR = 2(u_1 - u_2)/(u_1 + u_2)$ , and  $V_e = (u_1 + u_2)/2$ .

The instantaneous efficiency of a gear pair is defined as the ratio of the instantaneous output power to the input power

$$\eta(\phi_m) = \frac{L_{out}(\phi_m)\omega_{out}}{L_{in}\omega_{in}} \quad (29)$$

where  $L_{out}(\phi_m)$  and  $L_{in}$  are the values of torque acting on the output and input gears and  $\omega_{out}$  and  $\omega_{in}$  are the output and input rotational speeds, respectively. When only frictional losses are considered, the instantaneous output power can be written as the difference between the input power and the frictional power losses and hence the efficiency can be expressed as

$$\eta(\phi_m) = 1 - \frac{1}{L_{in}\omega_{in}} \sum_{q=1}^Q [ |F_s(u_1 - u_2)| + |F_r(u_1 + u_2)| ]_q \quad (30)$$

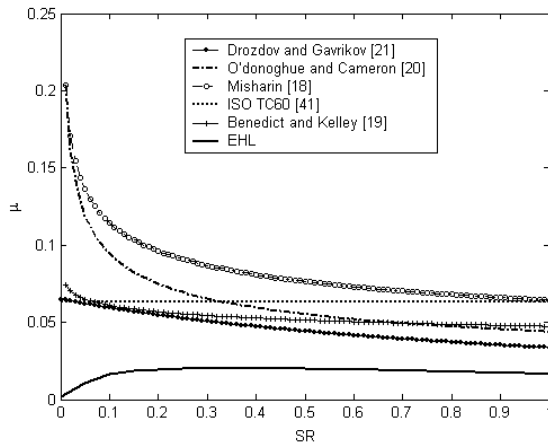
where the summation from  $q=1$  to  $Q$  is the sum of sliding frictional and rolling frictional power losses at position  $\phi_m$  and  $Q$  is the total number of load grids at the same mesh position.

## 2.5 Validation of Friction Coefficient Models

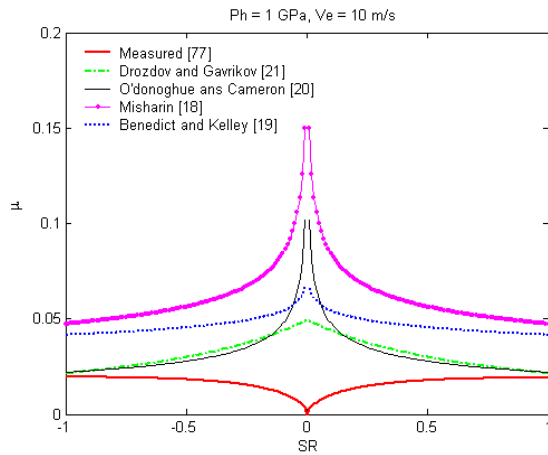
Comparison of Empirical and the EHL-Based  $\mu$  Models. In order to facilitate numerical comparisons among these formulae and the EHL model predictions, an example line contact problem with the following parameters is considered:

$$R = 5 \text{ mm}, V_e = 5 \text{ m/s}, P_h = 2 \text{ GPa}, \\ v_0 = 10 \text{ cPs}, v_k = 13 \text{ cSt}, S_{rms} = 0.07 \text{ } \mu\text{m} .$$

With these parameters, the calculated  $\mu$  values from published empirical formulae in Table 1, and



**Figure 3.** Comparison of friction coefficient models for  $R = 5 \text{ mm}$ ,  $V_e = 5 \text{ m/s}$ ,  $P_h = 2 \text{ GPa}$ ,  $\nu_0 = 10 \text{ cPs}$ ,  $\nu_k = 13 \text{ cSt}$ , and  $S_{rms} = 0.07 \text{ } \mu\text{m}$ .



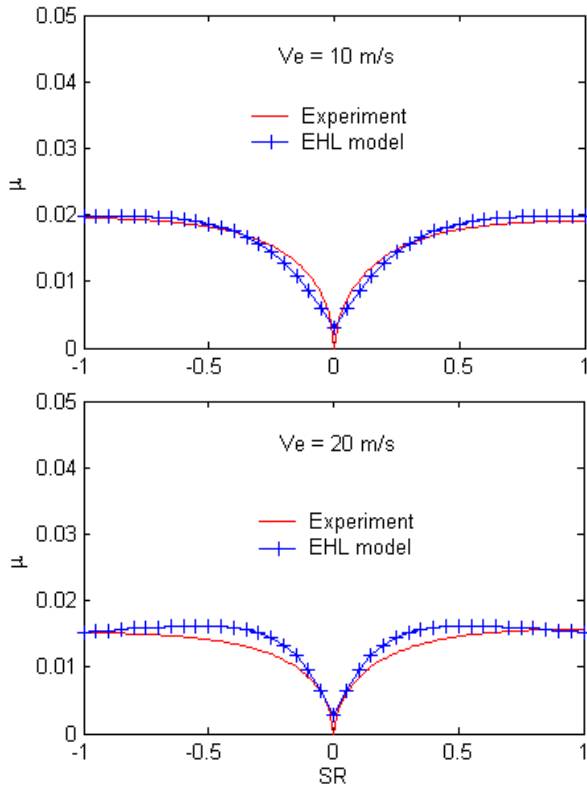
**Figure 4.** Comparison between published empirical friction coefficient formulae and the measured traction data.

from the thermal EHL model described in section 2.3 are plotted as a function of  $SR$  in Fig. 3. Here, three main discrepancies between the published empirical formulae and the thermal EHL-based  $\mu$  are evident. (i) When  $SR = 0$ , EHL model predicts nearly zero friction, while published empirical formulae all predict their largest  $\mu$  values. This is because  $V_s$  is in the denominator of most of these formulae. This implies that the friction coefficient is largest at the pitch point where there is no relative sliding in total conflict with the physical intuition and the EHL model prediction. (ii) The overall qualitative shape of the  $\mu$  versus  $SR$  curve from the EHL model is very different from those obtained from the other  $\mu$  formulae. (iii) The thermal EHL model predicts much lower  $\mu$  values than the empirical formulae, regardless of the value of  $SR$ . Since the differences amongst the individual  $\mu$  formulae as well as between the EHL model and these formulae are very significant, neither can be

used in the efficiency model with confidence since the resultant efficiency predictions will differ greatly. Only way to describe these discrepancies is to compare the predicted coefficients of friction to the measured ones.

A comparison between the published empirical  $\mu$  formulae and the measured traction data is shown in Fig. 4. The discrepancies between these published formulae and the fitted measured data are quite similar to those between the same formulae and the EHL model predictions in Fig. 3. In Fig. 4, for  $V_e = 10 \text{ m/s}$  and  $P_h = 1.0 \text{ GPa}$ , the measured and predicted  $\mu$  values differ in both magnitudes and qualitative shape. The published empirical formulae of others give larger  $\mu$  values regardless the value of  $SR$ . The predicted  $\mu$  values become very large as  $SR$  approaches zero while the measured  $\mu$  value is almost zero at  $SR = 0$ . Typical gear pairs might have  $SR$  values between  $-1$  and  $1$ . The same differences were observed at other load and speed conditions as well. The efficiency values predicted by using these published  $\mu$  formulae would be much lower than the actual mechanical efficiency values. Therefore, it can be concluded here that these formulae should not be used for gear efficiency studies and the results of previous efficiency models that employed these formulae [34-37, 40] should be taken cautiously. On the other hand, EHL model predictions agree well with the measured friction coefficient data for typical ranges of speed, load, and sliding ratio. The  $\mu$  values predicted by the EHL model and the measured values represented by the fitted formula are compared in Fig. 5 for  $P_h = 1.0 \text{ GPa}$  and  $V_e = 10$  and  $20 \text{ m/s}$ .

Based on this comparison, the EHL-based model can be considered valid. This is a positive outcome since the EHL model is physics-based requiring no empirical parameters. However, it presents a major difficulty especially when used in a gear efficiency model. Each EHL analysis takes several minutes of computations (more than 2 minutes of CPU time on a 3.0 GHz PC) and a complete gear efficiency analysis might require a few hundred EHL analyses. Therefore, this approach is not practical for real-life engineering applications requiring many parameter studies. In order to overcome this disadvantage, a design of experiment (DOE) of the EHL model is performed. Results from DOE were processed by multiple linear regression analysis to obtain a new  $\mu$  formula based on the EHL approach to be used in the gear efficiency models, eliminating the difficulties of this approach in terms of the computational demand. The new formula has the form



**Figure 5.** Comparison of EHL model predictions and the measured data at  $P_h = 1$  GPa and various  $V_e$  values.

$$\mu = e^{f(SR, P_h, v_0, S)} P_h^{b_2} |SR|^{b_3} V_e^{b_6} v_0^{b_7} R^{b_8} \quad (31)$$

where

$$f(SR, P_h, v_0, S) = b_1 + b_4 |SR| P_h \log_{10}(v_0) + b_5 e^{-|SR| P_h \log_{10}(v_0)} + b_9 e^S$$

Constants  $b_1$  to  $b_9$  vary based on type of lubricant used. This formula includes all the key features of a gear contact, namely  $SR$ ,  $V_e$ ,  $P_h$ ,  $S$ ,  $R$ , and  $v_0$ . It can be directly used to calculate  $\mu$  instead of running extremely time-consuming EHL analyses in real time, while still maintaining acceptable accuracy. It also avoids numerical issues associated with any EHL model. The same equation was used earlier in a spur gear mechanical efficiency prediction model that was validated through high-speed spur gear efficiency measurements [77].

### 3 Results and Discussions

With all necessary parameters at each contact point calculated by using the formulations and the contact model presented in previous sections, instantaneous efficiency of a hypoid gear pair can be calculated by applying the methodology defined in Fig. 1. A face-hobbed hypoid gear pair borrowed from an automotive application will be used here as an example system to demonstrate the hypoid gear

efficiency methodology. The main gear blank dimensions of this example system are listed in Table 2. The pinion member is generated and the ring gear is non-generated. A typical gear oil 75W90 is used as the lubricant whose properties are listed in Table 3.

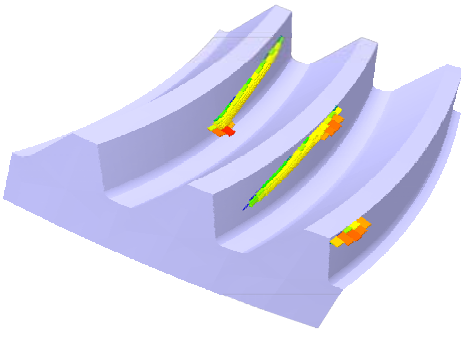
**Table 2** Example hypoid gear pair

Parameters	Pinion	Gear
Teeth number	12	41
Hand of spiral	Left	Right
Spiral angle (deg)	49.97	24.07
Outer cone distance (mm)	103.5	127.7
Face width (mm)	41.9	34.2
Face angle (deg)	28.6	58.7
Pitch angle (deg)	28.6	58.7
Front/Back angles (deg)	28.6	58.7
Pinion offset (mm)	44	
Shaft angle (deg)	90	
Material density (kg/m <sup>3</sup> )	7800	
Specific heat (J/kg K)	500	
Thermal conductivity (W/m K)	45	
Young's modulus (Pa)	2.07x10 <sup>11</sup>	
Poisson's ratio	0.3	

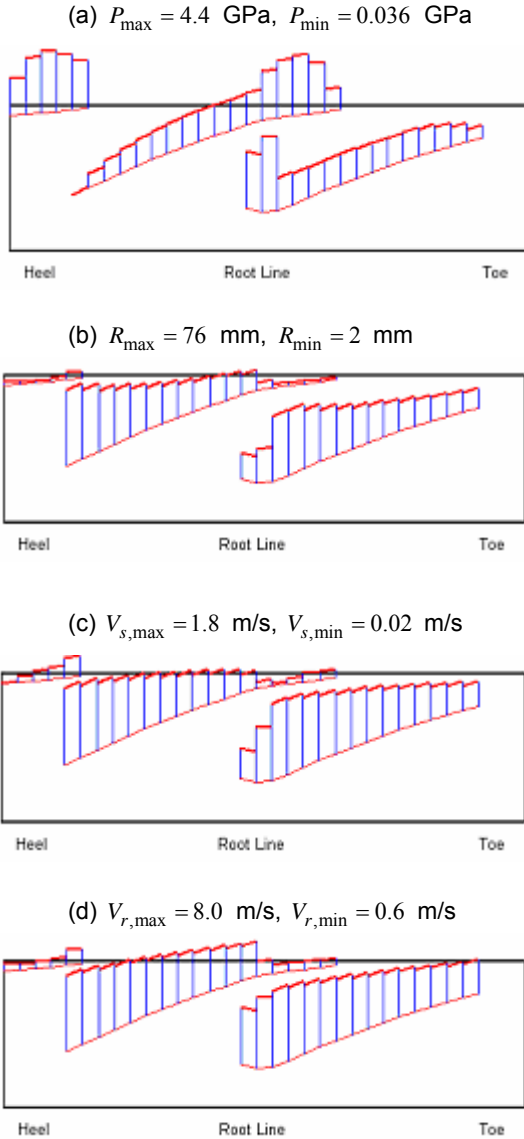
**Table 3** Lubricant parameters

Inlet temperature (K)	333
Viscosity (Pa s)	0.036
Density (kg/m <sup>3</sup> )	815
Pressure viscosity coefficient (1/Pa)	1.344x10 <sup>-8</sup>
Temperature viscosity coefficient	0.0217
Thermal conductivity (W/m K)	0.132
Coefficient of thermal expansion	6.5x10 <sup>-4</sup>
Specific heat (J/kg K)	2000

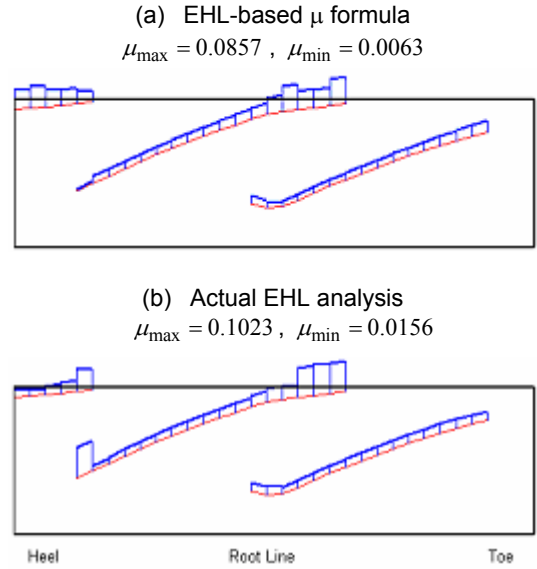
Figure 6 illustrates the contact pressures distribution at all the contact zones on the hypoid ring gear in contact at a given mesh position. As discussed earlier, contact pressures, radii of curvatures, sliding velocities, and sum of rolling velocities are four of the key geometry-related parameters used in the friction model. The distributions of these parameters along the contact lines at the same mesh angle  $\phi = \phi_1$  are shown in Fig. 7. It is clear from this figure that all four key parameters vary significantly along the contact lines.



**Figure 6.** Contact pressure distribution prediction of CALYX at a given mesh position  $\phi_1$ .



**Figure 7.** Distribution of (a) contact pressure, (b) combined radii of curvature, (c) sliding velocity, and (d) sum of rolling velocities along the normal of the instant line of contact at mesh position  $\phi_1$ .  $S = 0.1 \mu m$ ,  $L_{in} = 1,600 Nm$ ,  $N_p = 1,000 rpm$ , and  $T_{oil} = 60^\circ C$ .



**Figure 8.** Distribution of  $\mu(z, \theta, \phi_1)$  at every mesh grid of principal contact points predicted by using (a) the EHL-based  $\mu$  formula, and (b) actual EHL analysis.  $S = 0.1 \mu m$ ,  $L_{in} = 1,600 Nm$ ,  $N_p = 1,000 rpm$ , and  $T_{oil} = 60^\circ C$ .

The values of  $\mu(z, \theta, \phi_1)$  at each of the principal contact points along the instantaneous lines of contact are calculated by using the new formula in Eq. (31) as well as the actual EHL model are shown in Fig. 8 for  $S = 0.1 \mu m$ ,  $L_{in} = 1,600 Nm$ ,  $N_p = 1,000 rpm$ , and  $T_{oil} = 60^\circ C$ . For the EHL model prediction, a measured rough surface profile with  $RMS = 0.1 \mu m$  like the one shown in Fig. 9 is used. As shown in Fig. 8, the predicted  $\mu$  values show certain amount of variations along the instantaneous lines of contact. The  $\mu$  values become especially large when the contact occurs near the edge where contact pressures are large as shown in Fig. 7. The variations of  $\mu$  predicted at each contact point, together with the variations of the contact pressures and the sliding and rolling velocities, result in the variations of the power losses at each of these locations.

Figure 10 shows the variation of the average value of the instantaneous coefficient of friction at every mesh angle,  $\bar{\mu}(\phi_m) = Average[\mu(z, \theta, \phi_m)]$ , and the corresponding instantaneous mechanical efficiency  $\eta(\phi_m)$  of the gear pair. It is evident that the two EHL related models are in reasonably good agreement and give an average efficiency of about  $\eta = 98$  percent for  $S = 0.1 \mu m$ ,  $L_{in} = 1,600 Nm$ ,  $N_p = 1,000 rpm$  and  $T_{oil} = 60^\circ C$ .

A parametric study on the mechanical efficiency of the example hypoid gear pair is performed next. Figure 11 shows the variation of  $\bar{\eta}$  as a function of  $N_p$  for  $L_{in} = 500$  and  $1,000$ , given  $S = 0.4 \mu m$  and  $T_{oil} = 60^\circ C$ .  $L_{in}$  has a negligible influence on  $\bar{\eta}$

within this load range and curves for  $L_{in}=500$  and 1,000 Nm are almost identical. However,  $\bar{\eta}$  increases quite significantly with an increase in  $N_p$ . The mechanical efficiency values at  $N_p=500$  and 2,000 rpm differ by nearly 2 percent, which is due to very small rolling velocities at low speed that cause the friction coefficient increase significantly.

The combined influence of the surface roughness parameter  $S$  and the oil inlet temperature  $T_{oil}$  is illustrated in Fig. 12 for  $L_{in}=1,600$  Nm and  $N_p=1,000$  rpm. Both  $T_{oil}$  and  $S$  have a significant influence on  $\bar{\eta}$ . For instance at  $T_{oil}=60^\circ C$ ,  $\bar{\eta}=98.3$  percent for a smooth surface while it is 96.3 percent for  $S=0.6 \mu m$ . At the same roughness value of  $S=0.6 \mu m$ , elevating lubricant temperature to  $100^\circ C$  improves the efficiency from 96.3 percent to nearly 97.5 percent, partly due to the reduction in viscosity due to temperature increment [77,78].

Influence of hypoid gear assembly tolerances or errors,  $H$ ,  $V$ ,  $R$ , and  $\beta$  that are defined in Fig. 13 on  $\bar{\eta}$  are shown in Fig. 14. In Fig. 15,  $H$ ,  $V$ , and  $R$  are all varied from  $-0.5$  mm to  $+0.5$  mm and  $\beta$  is varied from  $-0.1$  degree to  $+0.1$  degree. From Fig. 14, it is found that shaft angle error  $\beta$  has almost no influence on  $\bar{\eta}$ , while influence of other three errors on  $\bar{\eta}$  is very limited for the cases considered.

#### 4 Conclusions

A model for the prediction of hypoid gears has been proposed and applied to a face-hobbed hypoid gear pair as a representative example. Formulations for the required surface velocities and radii of curvature have been presented. Influence of operating conditions, surface finish and lubricant temperature, as well as assembly errors on the efficiency of hypoid gears have also been quantified. While hypoid gear pair is defined by a very large number of parameters, extensive generic analyses were not possible to obtain information about the influence of each parameter on the mechanical efficiency of the hypoid gear pair, still several consistent trends were observed. The rotational speed, surface roughness amplitude, and lubricant temperature are all very influential on the mechanical efficiency of a hypoid gear pair. Meanwhile, the parameters such as load and assemble errors have a limited or no influence on hypoid efficiency.

#### Acknowledgement

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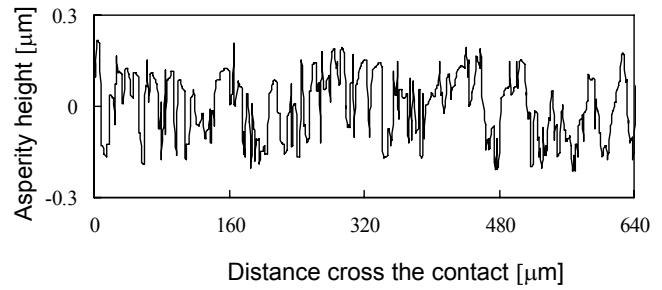


Figure 9. A measured surface roughness profile used in the EHL analysis.

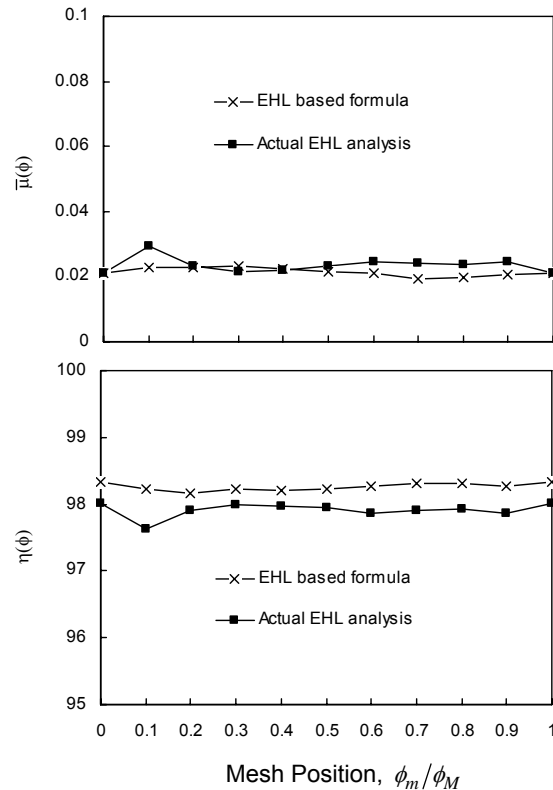
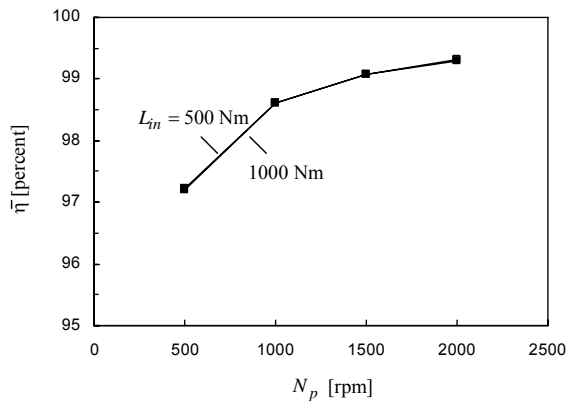
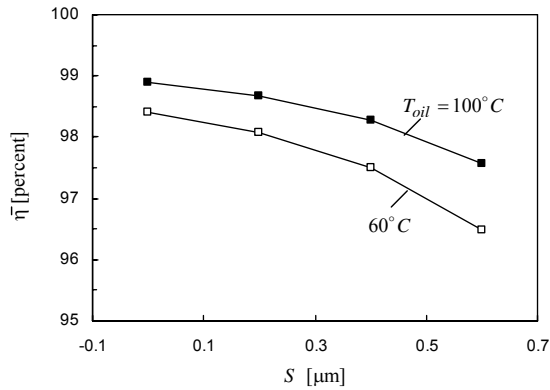


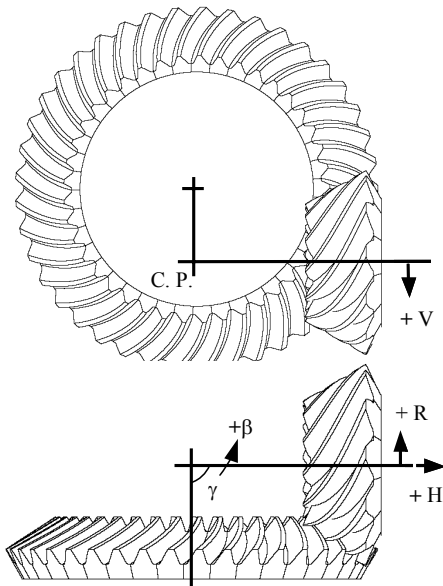
Figure 10.  $\bar{\mu}(\phi)$  and  $\eta(\phi)$  predicted by using the friction coefficient formulae developed in this study and the EHL model.  $S = 0.1 \mu m$ ,  $L_{in} = 1,600$  Nm,  $N_p = 1,000$  rpm, and  $T_{oil} = 60^\circ C$ .



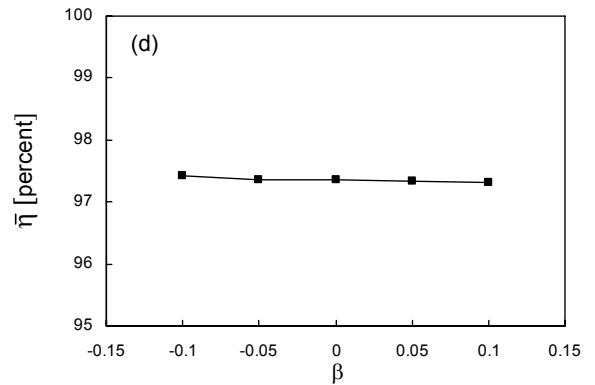
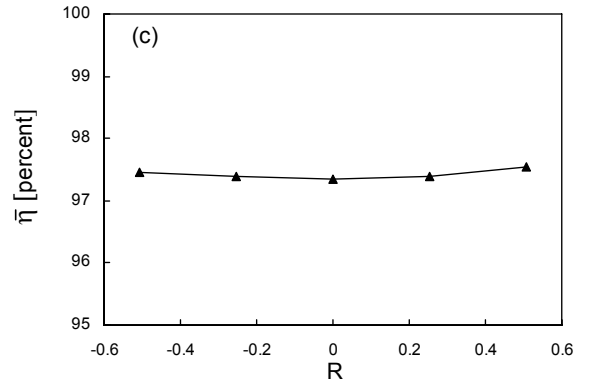
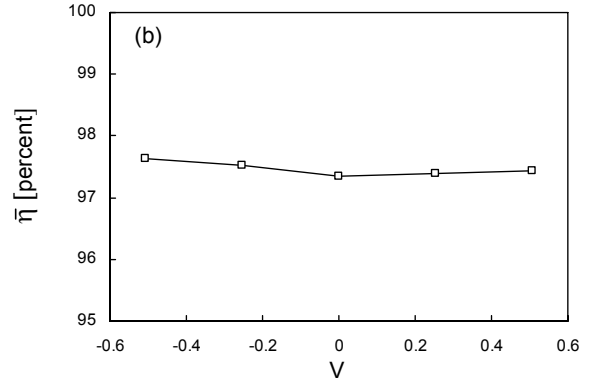
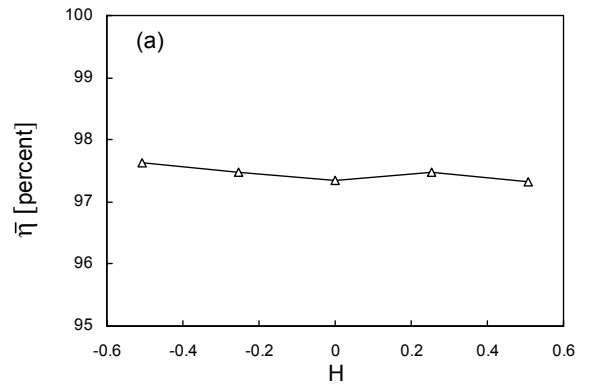
**Figure 11.** Influence of  $L_{in}$  and  $N_p$  on  $\bar{\eta}$ ;  $S = 0.4 \mu m$  and  $T_{oil} = 60^\circ C$ .



**Figure 12.** Influence of  $S$  and  $T_{oil}$  on  $\bar{\eta}$ ;  $L_{in} = 1,600 Nm$  and  $N_p = 1,000 rpm$ .



**Figure 13.** Example hypoid gear pair geometry and illustration of errors, V: pinion movement along offset, H: pinion axial movement, R: pinion movement along gear axis,  $\beta$ : shaft angle error.



**Figure 14.** Influence of errors (a) H, (b) V, (c) R, and (d)  $\beta$  on  $\bar{\eta}$ .  $S = 0.4 \mu m$ ,  $L_{in} = 1,600 Nm$ ,  $N_p = 1,000 rpm$ , and  $T_{oil} = 60^\circ C$ .

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